

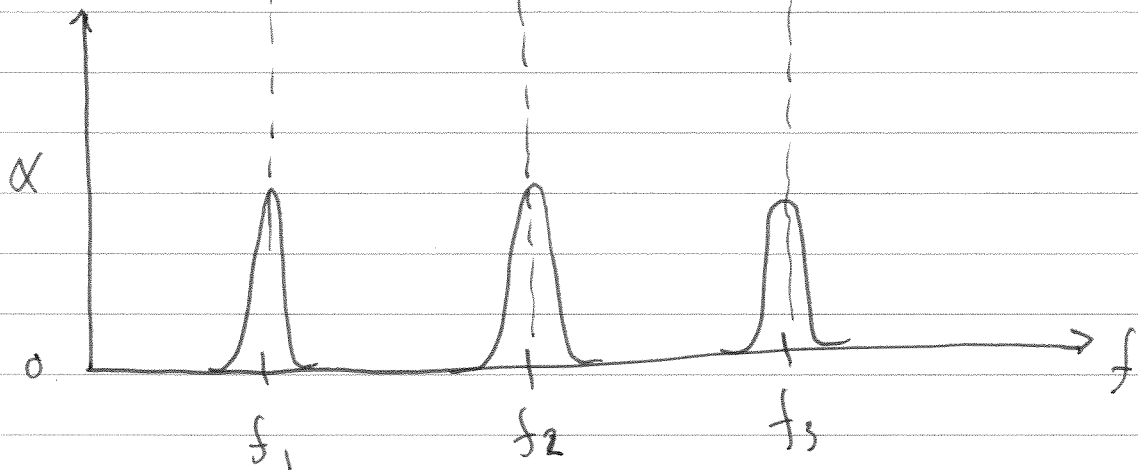
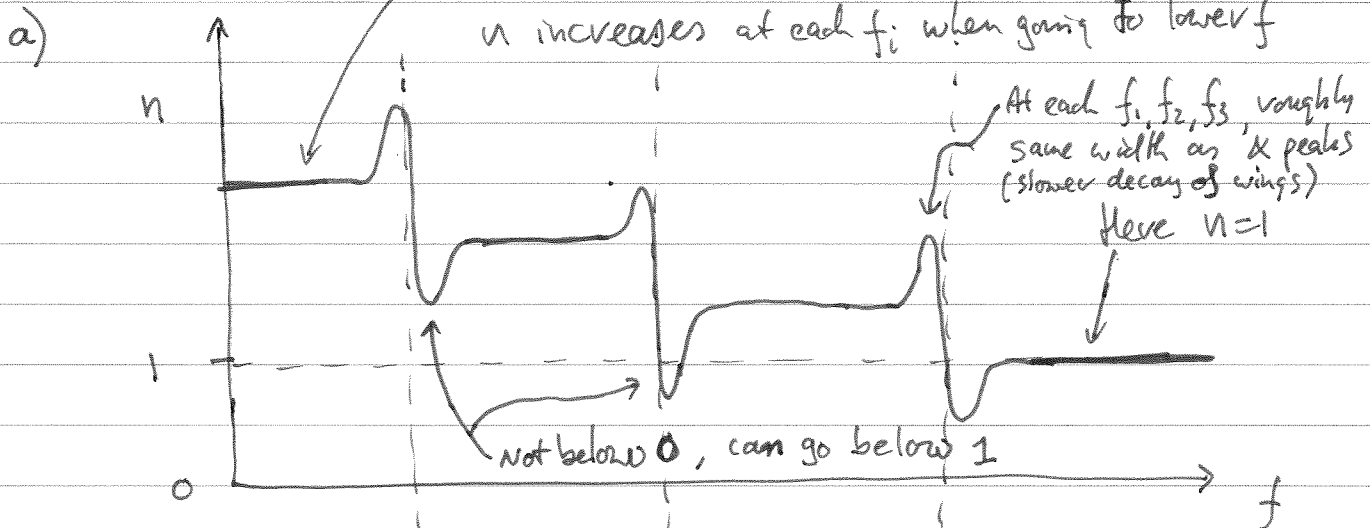
Model Answers Exam

Functional Properties - Optical Materials

26 Jan 2018

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Problem 1-opt



- b) The resonances f_2 and f_3 are at visible and UV frequencies. These typically correspond to electronic transitions. The resonance at f_1 is at infrared frequency, and could thus very well be a vibronic transition (lattice phonons).

c) It is not a metal or semiconductor, since there is never a wide band with strong absorption. It is probably not a molecular material, since these mostly have broader resonances and also broad phonon (vibronic) side bands (or phonon replicas) for f_2 and f_3 .

It could very well be a glass or crystalline insulator. Then f_1 is a distinct vibronic mode, and f_2, f_3 correspond to electronic transitions of localized impurity states.

f_3 is not an electronic transition of bulk material, since there is no increasing absorption for $f > f_3$. Apparently, this material is mostly transparent for $f > f_3$. That must be a material with a very large bandgap (indirect), and little X-ray absorption (light nuclei).

d) Estimate n for $f_n = 3 \times 10^{15} \text{ Hz} \Rightarrow$ is between f_2 and f_3 .
 \Rightarrow Only increase of n above 1 due to the resonance at f_3 .

Well below f_3 , we can use the approximation

$$\Delta E = E(f_n) - E(\infty) = E(f_n) - 1 = \frac{N_3 e^2}{\epsilon_0 m_0 \omega_0^2}$$

\Rightarrow Use here $\omega_0 = 2\pi f_3$
 m_0 is electron mass

$$n \approx \sqrt{E(f_n)} = \sqrt{1 + \Delta E}$$

$$= \sqrt{1 + \frac{N_3 e^2}{\epsilon_0 m_0 (2\pi f_3)^2}} \quad \begin{matrix} \text{fill in} \\ \text{numbers} \end{matrix} = 1.049$$

Problem 2-opt

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a) The value of R_x is the energy difference between the $n=1$ free exciton absorption line $\left[\left(E_{\text{gap}} - \frac{R_x}{n^2} \right) \text{ for } n=1 \right]$ and E_{gap} [that is, $\left(E_{\text{gap}} - \frac{R_x}{n^2} \right) \text{ for } n \rightarrow \infty$].

$$E_{\text{gap}} = E_g = 1.5192 \quad (\text{read from plot})$$

$$E_{x_{n=1}} = \underline{1.5149} \quad (\text{" " "})$$

$$R_x \approx 4.3 \text{ meV}$$

b) The Bohr radius of the free exciton is

$$r_n = \frac{m_0}{\mu} E_r n^2 a_H, \quad \text{where } a_H \text{ the Bohr radius}$$

of the Hydrogen atom for $n=1$, $a_H = 0.0529 \text{ nm}$.

Here μ is the reduced mass for the electron-hole system

$$\mu = \frac{1}{\frac{1}{m_e^*} + \frac{1}{m_h^*}} \quad \leftarrow \text{Electron and hole effective mass}$$

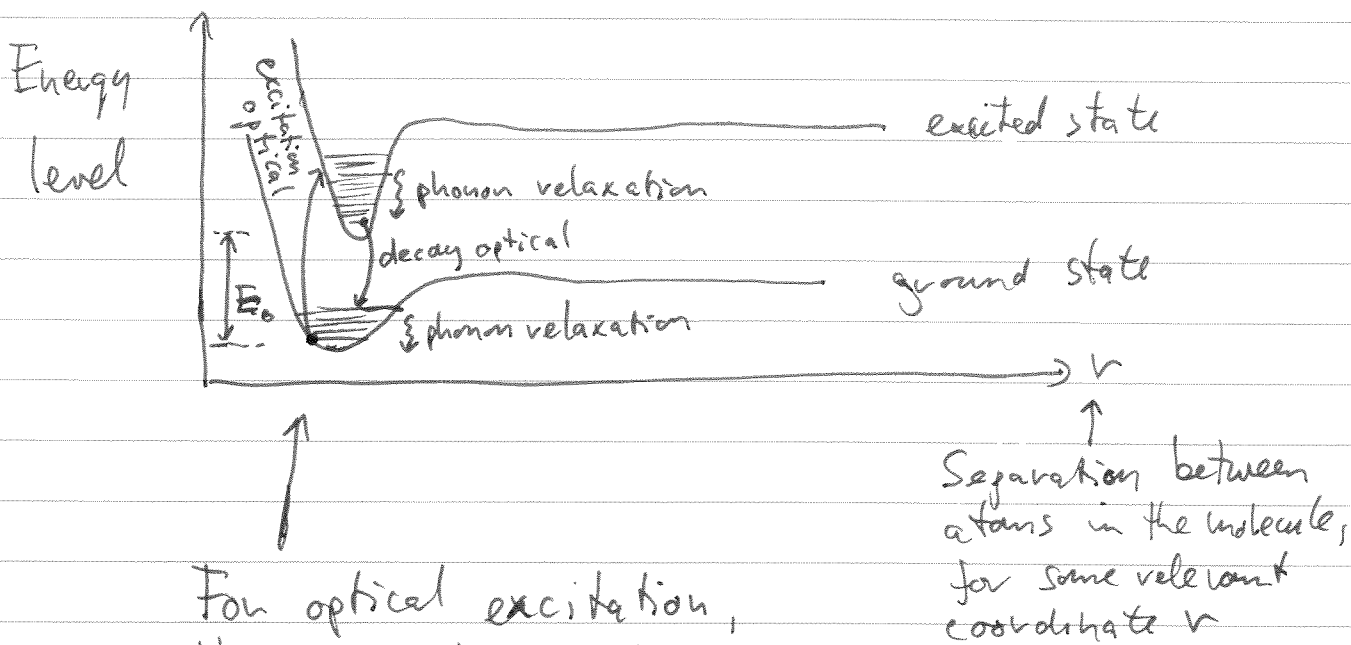
Filling in all numbers for r_n with $n=1$,
 $m_e^* = 0.067 m_0$, $m_h^* \approx 0.2 m_0$, $E_r = 13.6 \text{ eV}$

gives $r_n = 13.5 \text{ nm}$ \leftarrow estimate
(any value $m_e^* \leq m_h^* \leq m_0$ is graded as correct)

Problem 3-opt

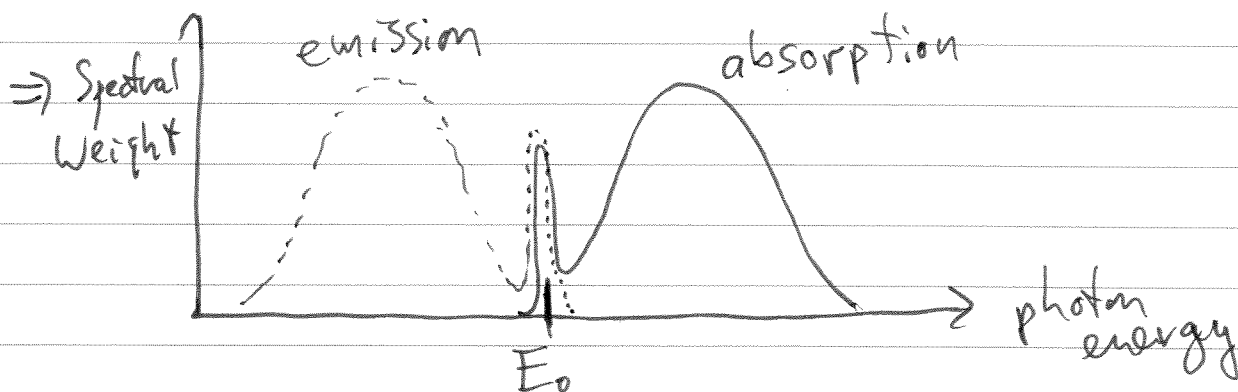
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Molecules have strong electronic and vibronic resonances that are coupled. This can be represented in the following diagram



For optical excitation, there is strong absorption for photon energies $> E_0$, due to a process that generates both the electronic transition and phonons

For optical emission, a big part of the photons come out at energies $< E_0$, since it can decay to the electronic ground state to a level with high phonon excitation.



Problem 4-opt

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a)

Potential	Curve	Explanation
B	2	This is the most parabolic curve, so the most linear response for the P-E relation
A	1	For this potential, the charged particles feel less restoring force at large displacement \Rightarrow P increases super-linear at high E driving amplitudes
C	3	For this potential, the charged particles feel an enhanced restoring force at large displacement \Rightarrow P increases sub-linear at high E driving amplitudes

b) $\chi^{\text{non linear}} = \chi^{(1)} + \chi^{(2)} E + \chi^{(3)} E^2$

$$P = \epsilon_0 \chi^{(1)} E + \epsilon_0 \chi^{(2)} E^2 + \epsilon_0 \chi^{(3)} E^3 + \dots$$

For each curve, the slope at $E=0$ is about the same and positive $\Rightarrow \chi^{(1)} > 0$ for all curves

For each curve, it is anti-symmetric about $E=0 \Rightarrow \chi^{(2)} = 0$ for all three curves

For curve 2, $\chi^{(3)} = 0$ since it is linear.

For curve 1, $\chi^{(3)} > 0$, since it is super-linear.

For curve 3, $\chi^{(3)} < 0$, since it is sub-linear.

c) If P oscillates at frequencies $\neq \omega$,

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this will cause the emission of EM fields at these different frequencies. For the case here, the response is dominated by $\chi^{(1)}$ and $\chi^{(3)}$.

Assume the driving field is $E = E_A \cos(\omega t) \Rightarrow$
↑ amplitude.

$$P = \epsilon_0 E_A \chi^{(1)} \cos(\omega t) + \epsilon_0 E_A^3 (\cos(\omega t))^3$$

Use $\cos^3(x) = \frac{3}{4} \cos(x) + \frac{1}{4} \cos(3x)$ \Rightarrow

$$P = \epsilon_0 E_A \chi^{(1)} \cos(\omega t) + \frac{3}{4} \epsilon_0 E_A^3 \chi^{(3)} \cos(\omega t) + \frac{1}{4} \epsilon_0 E_A^3 \cos(3\omega t)$$

P also oscillates at frequencies 3ω . The dominant frequencies in the response are ω and 3ω .